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Leibniz
Universität
Hannover



Fakultät für
Bauingenieurwesen
und Geodäsie



Internship in Saint Petersburg

September 18th, 2015

Mona Dannert

Supervisor: Vitaly Kuzkin

Content

- I) Introduction
- II) Fundamentals of Buckling
- III) Hoff's Problem
- IV) Analytical Solution
- V) Results
- VI) Conclusion

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I - Introduction

- Master Studies: Computational Engineering at Leibniz University, Hanover



<http://www.hannover.de/Wirtschaft-Wissenschaft/Wissenschaft/Hochschulen-und-Forschung/Hochschulen/>

I - Introduction

- Master Studies: Computational Engineering at Leibniz University, Hanover
- Structure of Studies
 - 2 semesters of lectures
 - 1 semester of internship (industry or abroad)
 - 1 semester Master's thesis

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- Master Studies: Computational Engineering at Leibniz University, Hanover
- Structure of Studies
 - 2 semesters of lectures
 - 1 semester of internship (industry or abroad)
 - 1 semester Master's thesis
- Abroad: get to know university research

I - Introduction

- Prof. Nackenhorst of Institute of Mechanics and Numerical Mechanics:
=> Department of Theoretical Mechanics

I - Introduction

- Prof. Nackenhorst of Institute of Mechanics and Numerical Mechanics:
 - => Department of Theoretical Mechanics
- Saint Petersburg:
 - arriving: April, 10th 2015
 - departure: October, 10th 2015
 - research on dynamic buckling of columns

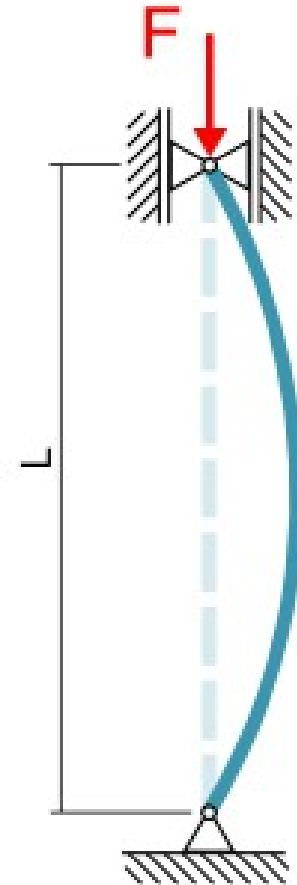
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II – Fundamentals of Buckling

- Statics

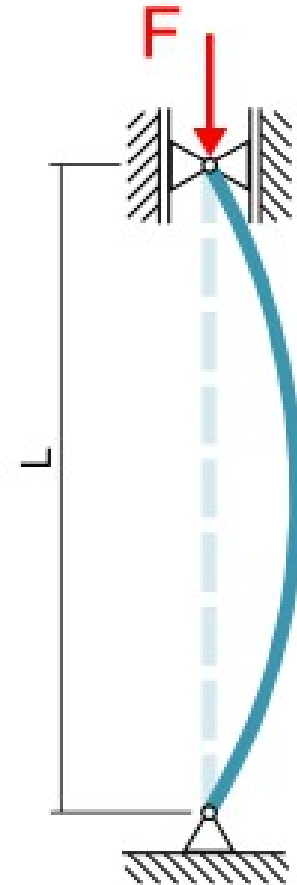
- Euler rod:
$$F_E = \frac{\pi^2 EI}{L^2}$$



<http://www.maschinenbau-wissen.de/skript3/mechanik/festigkeitslehre/134-knicken-euler>

II – Fundamentals of Buckling

- Statics
 - Euler rod: $F_E = \frac{\pi^2 EI}{L^2}$
- Dynamics
 - Buckling load exceeds Euler force



<http://www.maschinenbau-wissen.de/skript3/mechanik/festigkeitslehre/134-knicken-euler>

II – Fundamentals of Buckling

- Dynamics: three possibilities of applying load
 - suddenly applied force
 - moving one end by constant velocity
 - impact

II – Fundamentals of Buckling

- Dynamics: three possibilities of applying load
 - suddenly applied force
 - moving one end by constant velocity
 - impact

=> this is what happens in a common testing machine

Content

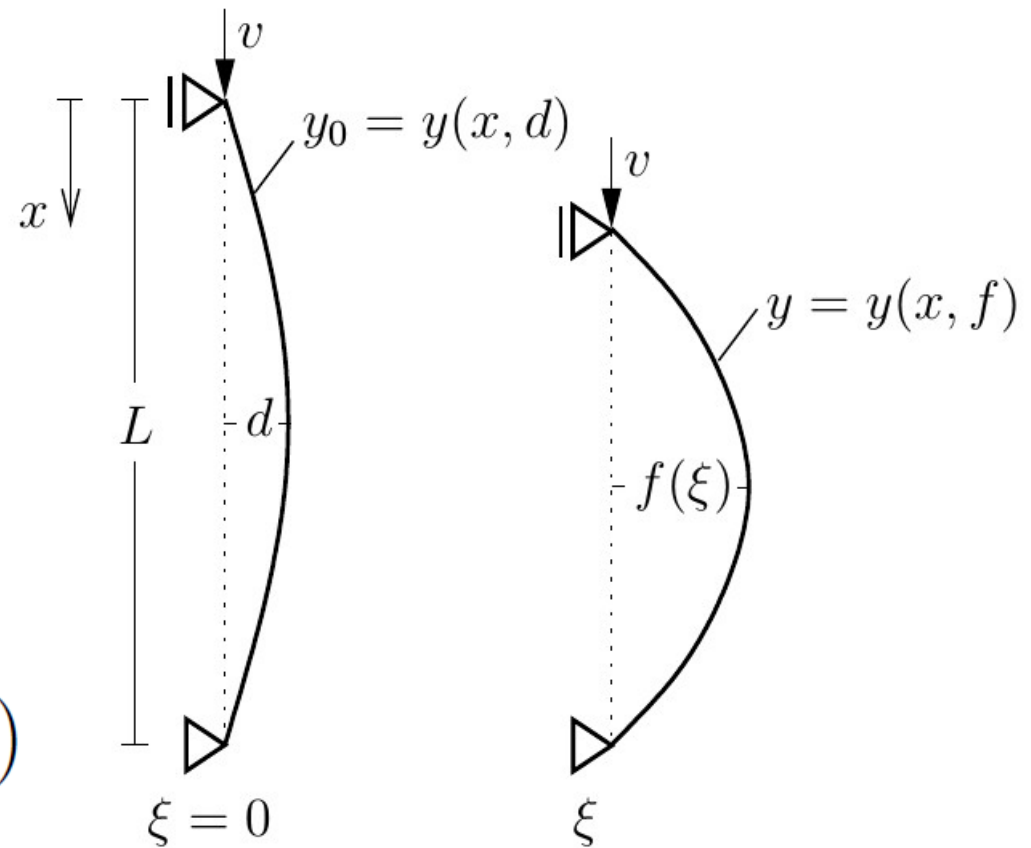
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III – Hoff's Problem

- Published in 1951
- Differential equation for dynamic buckling of an initially curved column
 - Describing a rod compressed by constant velocity
 - Influence of just two parameters

III – Hoff's Problem

- Initially curved column
 - initial deflection d
 - constant velocity v
 - deflection function $y = y(x, f)$
 - dimensionless amplitude $f = f(\xi)$



III – Hoff's Problem

- Transverse motion of the column

$$\frac{\partial^2}{\partial x^2} \left\{ EI \left[\frac{\partial^2 y}{\partial x^2} - \frac{\partial^2 y_0}{\partial x^2} \right] \right\} + P \frac{\partial^2 y}{\partial x^2} + \mu A \frac{\partial^2 y}{\partial t^2} = 0$$

- Deflection function $y = y(x, f) = Rf \sin \frac{\pi x}{L}$

III – Hoff's Problem

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- Deflection function $y = y(x, f) = Rf \sin \frac{\pi x}{L}$

- Boundary conditions $y(0, \xi) = y(L, \xi) = 0$

$$\frac{\partial^2 y}{\partial x^2} \Big|_{x=0} = \frac{\partial^2 y}{\partial x^2} \Big|_{x=L} = 0$$

III – Hoff's Problem

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- Deflection function $y = y(x, f) = Rf \sin \frac{\pi x}{L}$

- Compressing Force

$$P = \frac{EA}{L} \left\{ vt - \frac{1}{2} \int_0^L \left[\left(\frac{\partial y}{\partial x} \right)^2 - \left(\frac{\partial y_0}{\partial x} \right)^2 \right] dx \right\}$$

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III – Hoff's Problem

- Differential equation of dimensionless amplitude

$$f'' + \Omega \left[\frac{1}{4} f^3 + (1 - \xi) f - \frac{d^2}{4} f - d \right] = 0$$

- Boundary conditions $f(0) = d$ $f'(0) = 0$

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- Dimensionless time $\xi = \frac{vtL}{\pi^2 R^2}$

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- Dimensionless time $\xi = \frac{vtL}{\pi^2 R^2}$

- Similarity number $\Omega = \pi^8 \left(\frac{R}{L} \right)^6 \left(\frac{v_s}{v} \right)^2$

III – Hoff's Problem

- Differential equation of dimensionless amplitude

$$f'' + \Omega \left[\frac{1}{4}f^3 + (1 - \xi)f - \frac{d^2}{4}f - d \right] = 0$$

- Buckling time ξ_* by solving

$$\left. \frac{dP}{d\xi} \right|_{\xi=\xi_*} = EA \left(\frac{\pi R}{L} \right)^2 \left[1 - \frac{1}{2}f_*f'_* \right] = 0$$

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- Buckling force

$$\frac{P_*}{P_E} = \xi_* - \frac{1}{4}(f_*^2 - d^2)$$

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IV – Analytical Solution

- Assuming small imperfections
- Linearized differential equation

$$f'' + \Omega [(1 - \xi)f - d] = 0$$

IV – Analytical Solution

- Assuming small imperfections
- Linearized differential equation

$$f'' + \Omega [(1 - \xi)f - d] = 0$$

- Solution by using of Airy functions

$$f(\xi) = f_0 \pi \left[\text{Ai} \left(\Omega^{\frac{1}{3}}(\xi - 1) \right) \text{Bi}' \left(-\Omega^{\frac{1}{3}} \right) - \text{Bi} \left(\Omega^{\frac{1}{3}}(\xi - 1) \right) \text{Ai}' \left(-\Omega^{\frac{1}{3}} \right) \right] \\ - d \pi \Omega^{\frac{2}{3}} \left[\text{Ai} \left(\Omega^{\frac{1}{3}}(\xi - 1) \right) \int_0^{\xi} \text{Bi} \left(\Omega^{\frac{1}{3}}(z - 1) \right) dz \right. \\ \left. - \text{Bi} \left(\Omega^{\frac{1}{3}}(\xi - 1) \right) \int_0^{\xi} \text{Ai} \left(\Omega^{\frac{1}{3}}(z - 1) \right) dz \right].$$

IV – Analytical Solution

- Using asymptotic formula $\text{Ai}(z) \sim \frac{e^{-\frac{2}{3}z^{\frac{3}{2}}}}{2\sqrt{\pi}z^{\frac{1}{4}}}$
 - Assuming $z = \Omega^{\frac{1}{3}}(\xi - 1) > 0$ $\text{Ai} \rightarrow 0$

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$$- d \pi \Omega^{\frac{2}{3}} \left[\text{Ai}\left(\Omega^{\frac{1}{3}}(\xi - 1)\right) \int_0^{\xi} \text{Bi}\left(\Omega^{\frac{1}{3}}(z - 1)\right) dz \right.$$

$$\left. - \cancel{\text{Bi}\left(\Omega^{\frac{1}{3}}(\xi - 1)\right) \int_0^{\xi} \text{Ai}\left(\Omega^{\frac{1}{3}}(z - 1)\right) dz} \right].$$

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- With $f_0 = d$, solution can be simplified to
$$f(\xi) \approx \pi d \Omega^{\frac{1}{3}} \text{Bi} \left(\Omega^{\frac{1}{3}} (\xi - 1) \right)$$

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IV – Analytical Solution

- Solution can be simplified to

$$f_* = \sqrt{2} \left\{ \Omega \frac{3}{4} \ln \left[\frac{2}{\pi d^2 \Omega} \right] \right\}^{-\frac{1}{6}}$$

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=> restriction: $\Omega < \frac{4}{\sqrt{3}\pi d^2}$

IV – Analytical Solution

- Inserting solution

$$f_* = \sqrt{2} \left\{ \Omega \frac{3}{4} \ln \left[\frac{2}{\pi d^2 \Omega} \right] \right\}^{-\frac{1}{6}}$$

into $\frac{dP}{d\xi} \Big|_{\xi=\xi_*} = EA \left(\frac{\pi R}{L} \right)^2 \left[1 - \frac{1}{2} f_* f'_* \right] = 0$

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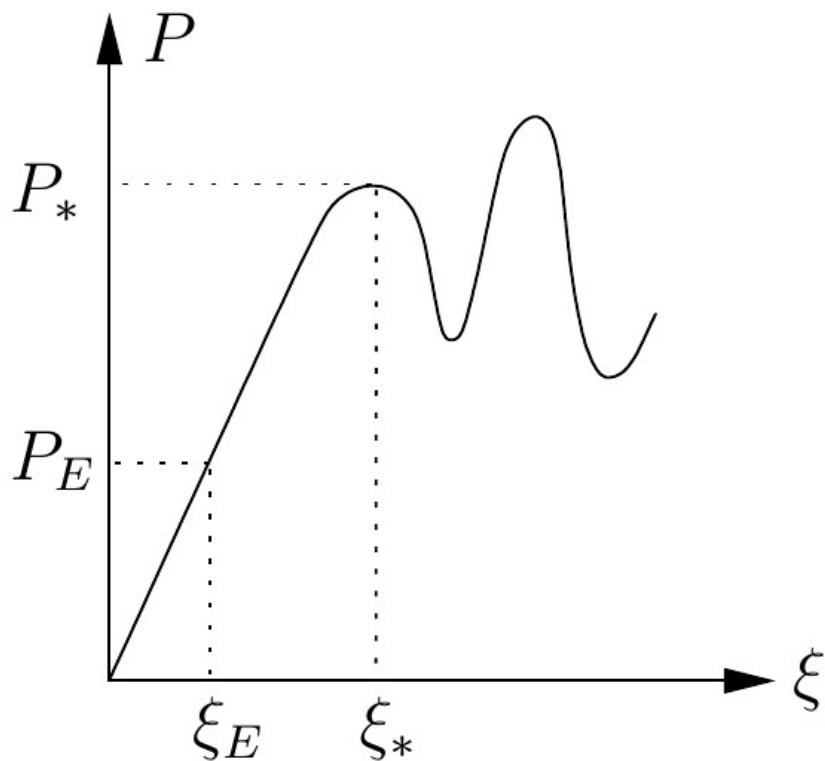
into $\left. \frac{dP}{d\xi} \right|_{\xi=\xi_*} = EA \left(\frac{\pi R}{L} \right)^2 \left[1 - \frac{1}{2} f_* f'_* \right] = 0$

=> Buckling time

$$\xi_* = 1 + \frac{1}{\Omega^{\frac{1}{3}}} \left\{ \frac{3}{4} \ln \left[\frac{2}{\pi d^2 \Omega} \right] \right\}^{\frac{2}{3}}$$

IV – Analytical Solution

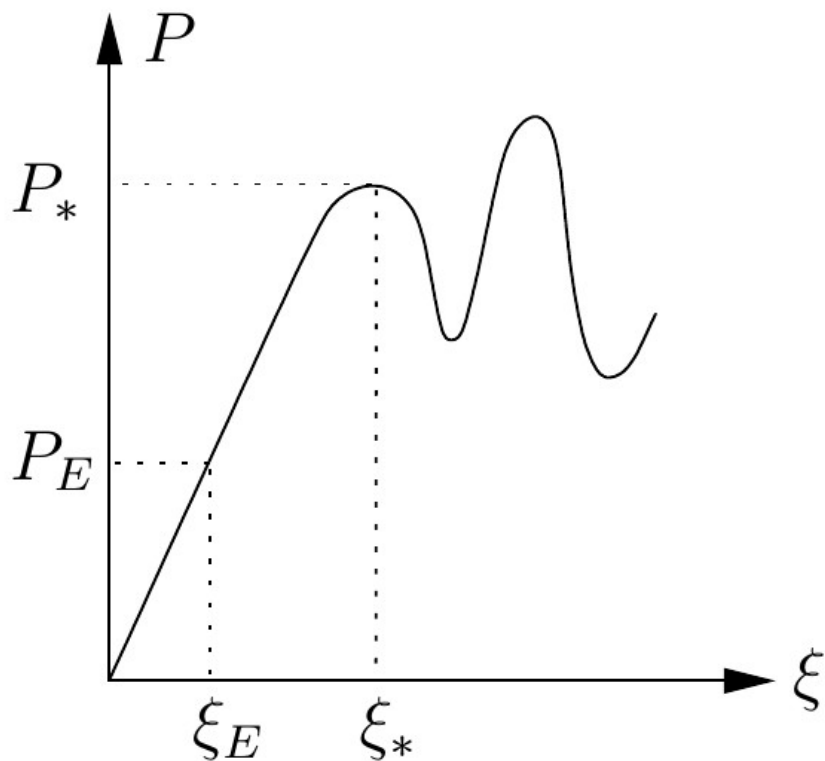
- Ratio of buckling force to Euler force



$$\frac{P_*}{P_E} = \xi_* - \frac{1}{4}(\cancel{f_*^2} - d^2)$$

IV – Analytical Solution

- Ratio of buckling force to Euler force

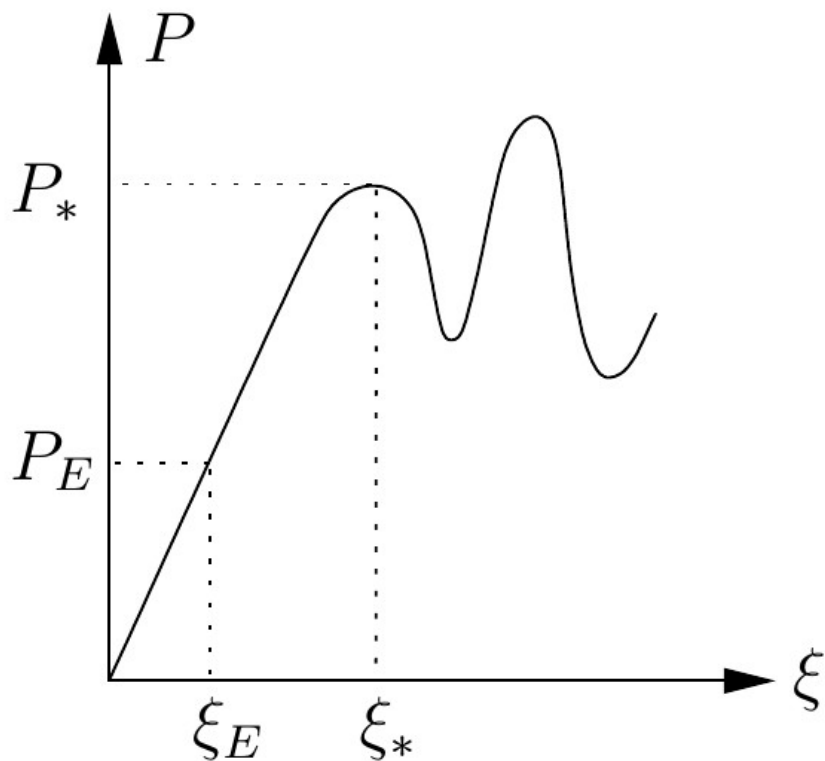


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$$\frac{P_*}{P_E} = 1 + \frac{1}{\Omega^{\frac{1}{3}}} \left\{ \frac{3}{4} \ln \left[\frac{2}{\pi d^2 \Omega} \right] \right\}^{\frac{2}{3}}$$

IV – Analytical Solution

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$$\frac{P_*}{P_E} = 1 + \frac{1}{\Omega^{\frac{1}{3}}} \left\{ \frac{3}{4} \ln \left[\frac{2}{\pi d^2 \Omega} \right] \right\}^{\frac{2}{3}}$$

- small initial deflection d

- similarity number $\Omega < \frac{4}{\sqrt{3}\pi d^2}$

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V – Results

- Computation with matlab

$$\frac{P_*}{P_E} = 1 + \frac{1}{\Omega^{\frac{1}{3}}} \left\{ \frac{3}{4} \ln \left[\frac{2}{\pi d^2 \Omega} \right] \right\}^{\frac{2}{3}}$$

V – Results

- Computation with matlab

$$\frac{P_*}{P_E} = 1 + \frac{1}{\Omega^{\frac{1}{3}}} \left\{ \frac{3}{4} \ln \left[\frac{2}{\pi d^2 \Omega} \right] \right\}^{\frac{2}{3}}$$

- assuming small initial deflection d
- following restriction $\Omega < \frac{4}{\sqrt{3}\pi d^2}$

V – Results

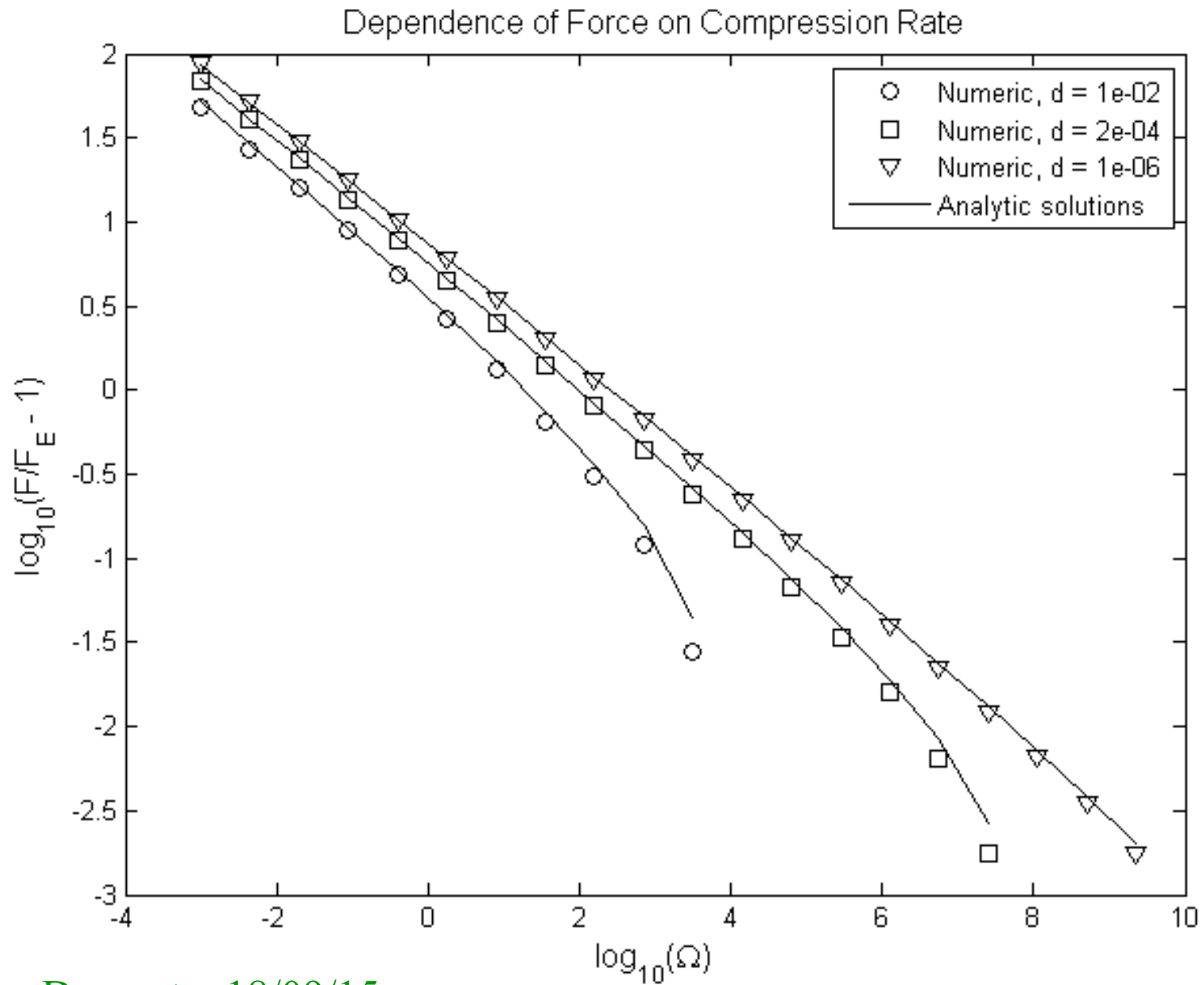
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- assuming small initial deflection d
- following restriction $\Omega < \frac{4}{\sqrt{3}\pi d^2}$

- Comparison with numerical solution of Hoff's differential equation

V – Results



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VI – Conclusion

- Solving Hoff's problem analytically
 - under certain conditions

VI – Conclusion

- Solving Hoff's problem analytically
 - under certain conditions
- Analytic solution approximates numeric one the better
 - the smaller initial deflection is
 - the smaller similarity number is

VI – Conclusion

- Solving Hoff's problem analytically
 - under certain conditions
- Analytic solution approximates numeric one the better
 - the smaller initial deflection is
 - the smaller similarity number is
- Dynamic buckling force tends to Euler force for large similarity numbers / small velocities

Literature

Hoff N.J. The dynamics of the buckling of elastic columns. *J. Appl. Mech.* 1951, Vol. 18, pp. 68–74.

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